A continuous-time version of the Derrida–Retaux model

Bastien Mallein Joint work with

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The Derrida–Retaux model

Let X_0 be a non-negative random variable. We set for all $n \in \mathbb{N}$,

$$X_{n+1} = (X_n^{(1)} + X_n^{(2)} - 1)_+,$$

where $X_n^{(1)}, X_n^{(2)}$ are two independent versions of X_n .

Question

What can be said of the asymptotic behavior of X_n as $n \to \infty$.

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Motivation

- Pinning models
- Parking in trees

2 A continuous-time version of this model

- Three intertwined models
- Results and open questions

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Poland–Scheraga model (1966)

Description

- DNA molecule of length N.
- Bounding energy ω_j in each site $j \leq N$.
- S_i distance between the two strands.





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- S_j distance between the two strands.

$$(S_i, 0 \le i \le N)$$
 path with $S_0 = S_N = 0$
 $(\omega_i, 0 \le i \le N)$ i.i.d.



Definition

Polymer measure of length N is defined by

$$\mathcal{P}_{N,\omega}(\mathrm{d}S) = rac{1}{Z_N} \exp\left(\sum_{i=1}^n \omega_i \mathbf{1}_{\{S_i=S_{i-1}=0\}}\right).$$

Remark

If $\omega_i > 0$, the strands are attracted to one another $\omega_i < 0$ they are repulsed.

Random polymer models, G. Giacomin, 2017.

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Définition

Set $F_{\omega} = \lim_{n \to \infty} N^{-1} \log Z_N$ the free energy of the model.

 $F_{\omega} = \log 2 \iff$ unpinned phase

$$\mathsf{F}_\omega < \mathsf{log}\, 2 \iff \mathsf{pinned} \mathsf{ phase}$$

Open problem

Find for which (ω_j) the model is in the pinned/unpinned phase. Find the behavior of the free energy around the pinned/unpinned phase transition.

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• Set $B \ge 2$ an integer.

- There is a unique edge at stage 0.
- Each edge is divided into *B* couple of edges when going from stage *k* to stage *k* + 1.
- Bounding energies are on the smallest edges.

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 ω_{\emptyset}

Figure: 0th stage

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Continuous Derrida–Retaux model





Figure: 2nd stage



Figure: 3rd stage



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Partition function

Set Z_k the sum of the energies of all possible paths at stage k, normalized by B^k . Note that the recursion equation implies

$$Z_{k+1} = rac{Z_k^{(1)}Z_k^{(2)} + B - 1}{B}.$$

The free energy is
$$F_{\omega} = \lim_{k \to \infty} \frac{1}{2^k} \log Z_k$$
.

Numerous results on that model : Monthus and Garet (2008), Derrida, Giacomin, Lacoin and Toninelli (2009), Lacoin and Toninelli (2009), Giacomin, Lacoin and Toninelli (2010, 2011), Berger and Toninelli (2013)...

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Continuous Derrida–Retaux mode

Derrida-Retaux model (2014)

Set $\bar{X}_k = \log Z_k$. The recursion equation can be rewritten

$$\bar{X}_{k+1} = \log(\exp(\bar{X}_k^{(1)} + \bar{X}_k^{(2)}) + B - 1)/B.$$

Derrida–Retaux model

Replace the function $x \mapsto \log((e^x + b - 1)/b)$ by $\max(x, -\log b)$.



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We obtain the recursion

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Free energy

Assume that $X_0 \sim p\delta_0 + (1-p)\mu$ with μ a probability distribution on $(0,\infty)$. Set

$$F_{\infty}(\rho) = \lim_{n \to \infty} 2^{-n} \mathbf{E}(X_n) \in [0, \infty],$$

and define $p_c \in [0,1]$ such that $F_{\infty}(p) > 0 \iff p < p_c$. We have $F_{\infty} = 0$ in the pinned phase, et $F_{\infty} > 0$ in the unpinned phase.

Conjecture

If $p_c>0$, we have $F_\infty(p)=\exp(-(c+o(1))(p-p_c)^{-1/2})$ as $p
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Assume that $X_0 \sim p\delta_0 + (1-p)\mu$ with μ a probability distribution on $(0,\infty)$. Set $F_{-}(p) = \lim_{n \to \infty} 2^{-n} F(X_n) \subset [0,\infty]$

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Nevertheless, it is reasonable to define Derrida-Retaux model on a tree.

Parking in a tree

Collet, Eckmann, Glaser, Martin (1984)

- Let X_0 be an integer-valued random variable.
- Consider a binary tree of height *n*.
- Put an independent copy of X₀ on each leaf of that tree.
- Then define recursively the value of each internal node to be $(a + b 1)_+$ if its children have value *a* and *b*.

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Theorem (Collet, Eckmann, Glaser, Martin (1984)) Recall that $F_{\infty} = \lim_{n \to \infty} 2^{-n} \mathbf{E}(X_n)$. We have

$$F_{\infty}=0\iff \mathbf{E}(X_02^{X_0})\leq E(2^{X_0})<\infty.$$

Moreover, if $F_{\infty} = 0$ then $\lim_{n\to\infty} X_n = 0$ in probability and if $F_{\infty} > 0$ then $\lim_{n\to\infty} 2^{-n}X_n = F_{\infty}$ in probability.

Open questions

If X_0 is not integer-valued, can we find a similar criterion distinguishing the pinned and unpinned phases ? Can we show the Derrida–Retaux conjecture on the behavior of the free energy around the pinning transition ?

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Recall that $F_{\infty}(p)$ is the free energy of $X_0 \sim p\delta_0 + (1-p)\mu$.

Theorem (Chen, Dagard, Derrida, Hu, Lifshits, Shi 2019) If μ is integer-valued and $\int x^3 2^{x} \mu(dx) < \infty$, then

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In their article, Derrida and Retaux write that a "good" renormalisation of the DR model should converge in law towards a process $(X_t, t \ge 0)$ whose law $\tilde{\varrho}_t$ satisfies the partial differential equation

$$\forall x > 0, \partial_t \tilde{\varrho}(x) = \partial_x \tilde{\varrho}(x) + \tilde{\varrho} * \tilde{\varrho}(x) - \tilde{\varrho}(x).$$

Problem

This equation is not conservative : even if ρ_0 is a probability measure, the total mass of ρ_t will be different from 1 for t > 0 large enough.

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Aim

Construct a stochastic process whose density satisfies a PDE close to the one given by Derrida and Retaux.

Remark

Observe that a process version of the Derrida–Retaux model can written as follows: given (U_n) i.i.d. variables with uniform distributions on [0, 1], we set

$$\begin{cases} X_{n+1} = X_n + F_n^{-1}(U_{n+1}) - \mathbf{1}_{\{X_n + F_n^{-1}(U_{n+1}) > 0\}} \\ F_{n+1}(x) = \mathbf{P}(X_{n+1} \le x) \text{ for all } x \ge 0. \end{cases}$$

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Definition

A continuous-time DR model is a solution of the following SDE:

$$\begin{cases} X_t = X_0 - \int_0^t \mathbf{1}_{\{X_s > 0\}} \mathrm{d}s + \int_0^t \int_0^1 F_s^{-1}(u) \mathcal{N}(\mathrm{d}s, \mathrm{d}u) \\ F_t(x) = \mathbf{P}(X_t \le x) \text{ for all } x \ge 0 \end{cases}$$

where N is a PPP on $\mathbb{R}_+ \times [0,1]$ with intensity $\mathrm{d}s\mathrm{d}u$.

Let us compute the equation satisfied by the law ρ_t of X_t . Let f be a C^1 function, we have

$$\mathbf{E}(f(X_t)) = \mathbf{E}(f(X_0)) - \int_0^t \mathbf{E}(f'(X_s)\mathbf{1}_{\{X_s>0\}}) \mathrm{d}s$$

+ $\int_0^t \mathbf{E}(f(X_s + X'_s) - f(X_s)) \mathrm{d}s$
$$\int f(x) \mathrm{d}\varrho_t(x) = \int f \mathrm{d}\varrho_0(x) - \int_0^t \int f'(x)\mathbf{1}_{\{x>0\}} \mathrm{d}\varrho_s(x) \mathrm{d}s$$

+ $\int_0^t \int f(y) \mathrm{d}\varrho_s * \varrho_s(y) \mathrm{d}s - \int_0^t \int f(x) \mathrm{d}\varrho_s(x) \mathrm{d}s.$

As a result, we obtain the PDE

$$\partial_t \varrho = \partial_x \left(\mathbf{1}_{\{x > \mathbf{0}\}} \varrho \right) + \varrho * \varrho - \varrho.$$

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time↓

Figure: Solving the stochastic differential equation

$$X_t = X_0 - \int_0^t \mathbf{1}_{\{X_s > 0\}} \mathrm{d}s + \int_0^t \int_0^1 F_s^{-1}(u) N(\mathrm{d}s, \mathrm{d}u).$$

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Figure: Solving the stochastic differential equation

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• Take a **Yule tree** of height *t*.

- Start with i.i.d. amout of paint on each leaf of the tree, sampled according to the law of X₀.
- Paint down the branches with a quantity 1 of paint per unit of branch length, until no more paint is left.
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Connexion between the results

Proposition

There exists a unique weak solution to the McKean-Vlasov type SDE, with same law at time t as in the tree-painting scheme of height t.

Proposition

There exists a unique weak solution to the PDE

$$\partial_t \varrho = \partial_x (\mathbf{1}_{\{x>0\}} \varrho) + \varrho * \varrho - \varrho,$$

Moreover, ρ_t is the law of the solution of the SDE at time t.

Proof.

Using Itô calculus, the law of X_t is a solution to the SDE Reciprocally, we prove that any solution to this PDE can be rewritten in a tree-painting scheme.

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Proposition

There exists a unique weak solution to the PDE

$$\partial_t \varrho = \partial_x (\mathbf{1}_{\{x>0\}} \varrho) + \varrho * \varrho - \varrho,$$

Moreover, ρ_t is the law of the solution of the SDE at time t.

Proof.

Using Itô calculus, the law of X_t is a solution to the SDE Reciprocally, we prove that any solution to this PDE can be rewritten in a tree-painting scheme.

Proposition

- The free energy $F_{\infty} = \lim_{t \to \infty} e^{-t} \mathbf{E}(X_t)$ exists.
- If $F_{\infty} > 0$, then $\lim_{t\to\infty} e^{-t}X_t = F_{\infty} \operatorname{Exp}(1)$ in law.

Proof.

1. By the SDE representation of the DR model, we have

$$\mathbf{E}(X_t) = \mathbf{E}(X_0) - \int_0^t \mathbf{P}(X_s > 0) \mathrm{d}s + \int_0^t \mathbf{E}(X_s) \mathrm{d}s$$

Solving this ODE we have $\mathbf{E}(X_t) = e^t(\mathbf{E}(X_0) + \int_0^t e^{-s} \mathbf{P}(X_s > 0) ds)$, hence the limit exists. **2.** We use that for a Yule tree, $e^{-s} \# \mathcal{N}_s \to \text{Exp}(1)$ in law.

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Open questions

Conjecture

If $F_{\infty} = 0$, then $\lim_{t \to \infty} X_t = 0$ in probability.

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If $F_{\infty} = 0$, the law of X_t , conditioned on $X_t > 0$ converges in law to a standard exponential distribution.

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Plan

Motivation

- Pinning models
- Parking in trees
- 2 A continuous-time version of this model
 - Three intertwined models
 - Results and open questions

3 An exactly solvable version of this model

We look for solution of the PDE taking the form

$$\varrho_t(\mathrm{d} x) = p_t \delta_0(\mathrm{d} x) + (1 - p_t) \lambda_t e^{-\lambda_t x} \mathrm{d} x.$$

The PDE can then be rewritten as

$$\begin{cases} p' = (1 - p)(\lambda - p)\\ \lambda' = -\lambda(1 - p) \end{cases}$$

Property

Mixtures of exponential random variables are stable under the action of the PDE (interpreting Dirac mass at 0 as exponential random variable with infinite parameter).

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Consider a solution of the system of differential equations

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The quantity $H := \frac{p_t}{\lambda_t} - \log \lambda_t$ is a preserved quantity.

Observe as well that λ is non-increasing.

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Proposition

Let (p, λ) be a solution of the differential system.

• If (p_0, λ_0) in the supercritical phase, then

 $\lambda_t \sim Ke^{-t}$ and $p_t \sim Kte^{-t}$.

If (p_0, λ_0) in the subcritical phase phase, then for some x > 1,

$$\lambda_t - x \sim K e^{-t(x-1)}$$
 and $1 - p_t \sim K rac{x-1}{x} e^{-(x-1)t}$

If (p₀, λ₀) on the critical line, then
$$\lambda_t = 1 + \frac{2}{t} + \frac{8\log t(1+o(1))}{3t^2} \quad \text{and} \quad p_t = 1 - \frac{2}{t^2} + \frac{16\log t(1+o(1))}{3t^3}.$$

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The Derrida-Retaux conjecture

Recall that

$$F_{\infty} = \lim_{t \to \infty} e^{-t} \mathbf{E}(X_t) = \lim_{t \to \infty} e^{-t} \frac{1 - p_t}{\lambda_t}.$$

With careful analysis of the differential equation, one proves

Theorem

Fix $\lambda \in (0, e)$ and set $p_c = \lambda \log \lambda - \log \lambda$, we have

$$\begin{split} F_{\infty}(p,\lambda) &\sim C \exp(-\pi\sqrt{2\lambda}(p_c-p)^{-1/2}) \text{ as } p \uparrow p_c \text{ if } \lambda > 1\\ F_{\infty}(p,\lambda) &\sim C(1-p)^{2/3} \exp(-\frac{\pi}{\sqrt{2}}(1-p)^{-1/2}) \text{ as } p \uparrow 1 \text{ if } \lambda = 1\\ F_{\infty}(p,\lambda) &\sim C(1-p)^{1/(1-\lambda)} \text{ as } p \uparrow 1 \text{ if } \lambda < 1 \end{split}$$

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- On the event $X_t > 0$, $N_t = O(t^2)$ leaves will contribute to the mass of paint at the origin.
- The total mass M_t of paint that was put on these leaves is approximately c_*N_t .
- The tree of the origin of the paint scales towards a time-inhomogeneous branching Markov process: particles grow mass linearly, and a particle of mass m splits at rate $2m/(1-t)^2$.



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- Can the results given for exponential variables be extended to general initial measures ?
- ② Can the results be extended to the original Derrida–Retaux model ?
- Or a similar behavior be observed for a probability distribution satisfying

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Thank you for your attention!