

# Bifurcations through a non-semisimple hyperbolic 1:1 resonance, with examples in rigid body dynamics

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International Conference “Dynamics in Siberia”

02.03.2026

Novosibirsk

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- [2] A.Z. Ali, V.A. Kibkalo, E.A. Kudryavtseva, M.V. Onufrienko, Bifurcations in integrable Hamiltonian systems near corank-one singularities // *Diff. Eq.* **60**:10 (2024), 1311–1368.
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Hamiltonian system with  $n = 2$  degrees of freedom ( $\Omega = \sum_{j=1}^n dp_j \wedge dq_j$ ) depending on a parameter  $\varepsilon \in \mathbb{R}$ :

$$\frac{dq_i}{dt} = \frac{\partial H^\varepsilon}{\partial p_i}(q, p), \quad \frac{dp_i}{dt} = -\frac{\partial H^\varepsilon}{\partial q_i}(q, p), \quad i = 1, 2. \quad (1)$$

Suppose  $F^\varepsilon$  is a first integral, i.e.  $\{H^\varepsilon, F^\varepsilon\} = 0$  for each  $\varepsilon$ , where  $\{f, g\} = \Omega(X_f, X_g)$ ,  $dg = \omega(\cdot, X_g)$ .

Suppose that the origin is an equilibrium (a singular point of rank 0) at  $\varepsilon = 0$ :

$$H^0 = H_2^0 + H_3^0 + \dots, \quad F^0 = F_2^0 + F_3^0 + \dots$$

Пусть  $\{\pm\lambda_1, \dots, \pm\lambda_n\}$  — спектр гамильтонова оператора  $\Omega^{-1}d^2H(x_0)$ .

**Theorem. (Int) [Birkhoff], [Bryuno 1988]** If  $\nexists \mathbf{k} \in \mathbb{Z}^n \setminus \{0\}$ ,  $\langle \mathbf{k}, \vec{\lambda} \rangle = 0$ , then  $\exists$  local formal symplectic coordinates  $(P, Q)$ , in which  $H = h(B_1, \dots, B_n)$ ,

$$B_j = \begin{cases} \frac{1}{2}(P_j^2 + Q_j^2) & \text{(center),} \\ P_j Q_j & \text{(saddle),} \end{cases} \quad B_{j-1} + iB_j = (P_{j-1} + iP_j)(Q_{j-1} - iQ_j) \text{ (focus-focus),}$$

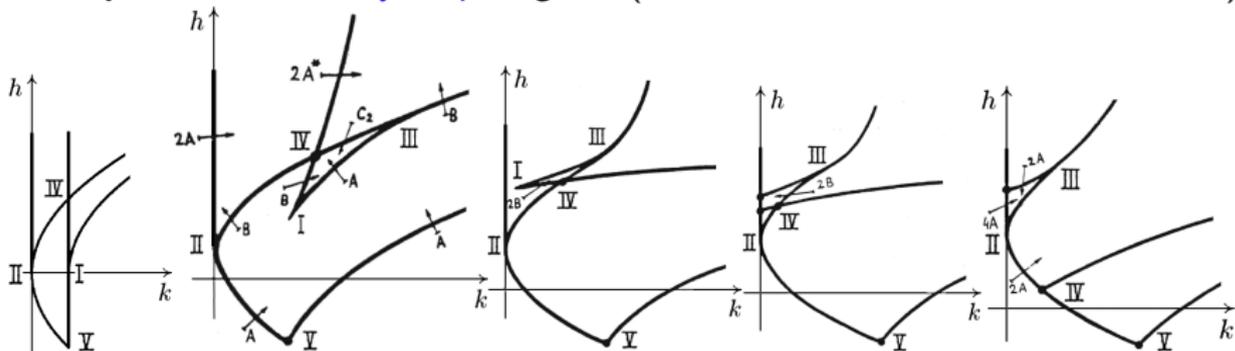
**(Stab)** if  $d^2H^0(0)$  is positive definite, then  $x_0$  is Lyapunov stable (elliptic); if  $\exists \lambda_j \notin i\mathbb{R}$ , then  $x_0$  is not Lyapunov stable,

**(Bifurc)**  $T_0\lambda_1 = 2\pi i$ ,  $\frac{\lambda_j}{\lambda_1} \notin \mathbb{Z} \Rightarrow \exists$  a family of periodic solutions  $N_1 \approx D^2(P_1, Q_1)$ ,

$H|_{N_1} = h(P_1^2 + Q_1^2)$  (**normal Lyapunov mode**;  $\frac{\lambda_1}{\lambda_2} = \frac{m}{n}$  [Duistermaat 1984]),

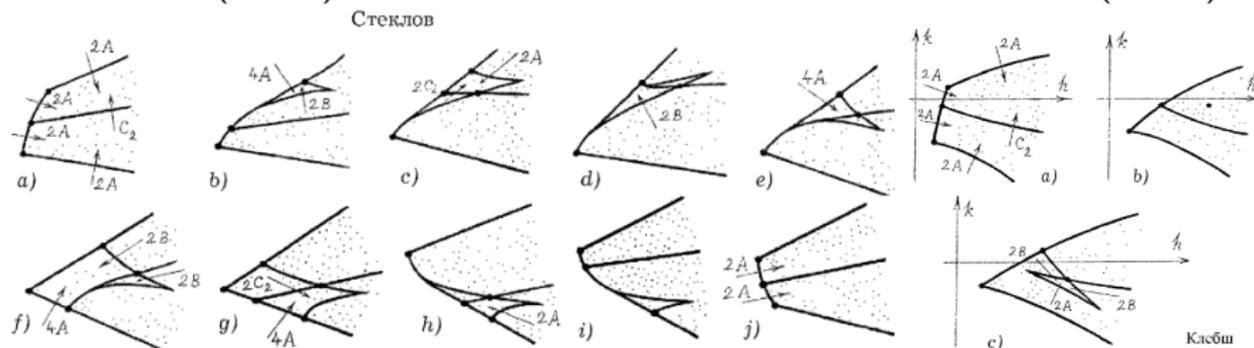
**(StrStab)**  $\forall$  small perturbation  $\tilde{H} \exists!$  equilibrium  $\tilde{x}_0 \in U$ .

**Definition** of **non-semisimple hyperbolic 1 : 1 resonance** at the equilibrium  $(q, p) = (0, 0)$ . Suppose that the Hamiltonian operators  $A = X_{H_2^0}$ ,  $X_{F_2^0}$  are **linearly independent**, and  $A$  is **non-diagonalizable** (Jordan block structure), so it has a multiple spectrum  $\text{spec } A = \{\lambda, \lambda, -\lambda, -\lambda\}$ , which we suppose to be **real**,  $\lambda > 0$ .  
**Examples:** **Kovalevskaya top** at  $g = 1$  ( $b \rightarrow c$ , two anti-canonical **involutions**):



**Steklov case** ( $d \rightarrow e$ ),

**Clebsch case** ( $b \rightarrow c$ ):



**Properties:** 1)  $\exists$  **linear changes** of variables  $(q, p) \rightarrow (\tilde{q}, \tilde{p})$  and  $(H, F) \rightarrow (H, \tilde{F} = aH + bF)$  transforming  $H_2^0, F_2^0$  to the form

$$H_2^0 = \lambda \tilde{S} + \tilde{X}, \quad \tilde{F}_2^0 = \tilde{S}, \quad \text{where } S = p_1 q_1 + p_2 q_2, \quad X = p_1 q_2, \quad Y = p_2 q_1, \quad Z = p_1 q_1 - p_2 q_2$$

1')  $\exists$  non-linear local changes  $(q, p) \rightarrow (\tilde{q}, \tilde{p}) = (\tilde{q}^\varepsilon(q, p), \tilde{p}^\varepsilon(q, p))$  and  $(H, F) \rightarrow (H, \tilde{F} = \tilde{F}^\varepsilon(H, F))$  transforming  $H_2^0, F_2^0$  to a **normal form**

$$H^\varepsilon = \lambda(\varepsilon) \tilde{S} + \tilde{X} + \varphi^\varepsilon(\tilde{Y}, \tilde{S}), \quad \tilde{F}^\varepsilon = \tilde{S},$$

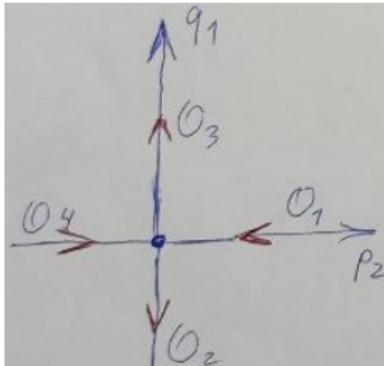
2) equilibria (near the semi-simple 1:1 resonance) form a **1-parameter family**  $(\tilde{q}, \tilde{p}) = (0, 0)$  with parameter  $\varepsilon$ , so  $\text{spec } A^\varepsilon = \{\pm \lambda_1^\varepsilon, \pm \lambda_2^\varepsilon\}$  where  $\lambda_1^\varepsilon$  and  $\lambda_2^\varepsilon$  collide at  $\varepsilon = 0$ ,

3) locally  $\exists$  two (families of commuting) **anti-canonical involutions**

$\tau_\pm^\varepsilon : (\tilde{q}_1, \tilde{q}_2, \tilde{p}_1, \tilde{p}_2) \mapsto \pm(\tilde{p}_2, \tilde{p}_1, \tilde{q}_2, \tilde{q}_1)$ , so the system is reversible w.r.t. them,

3') singular points on the integral manifold  $L^{0,0} := \{H^0 = F^0 = 0\}$  (at which  $dH$  and  $dF$  are linearly dependent) form (locally) four rays

$\{p_1 = q_2 = 0, p_2 q_1 = 0\} = 0 \cup \mathcal{O}_1 \cup \dots \cup \mathcal{O}_4$  where  $\mathcal{O}_j$  are 4 **1-dimensional local orbits** of the local Hamiltonian action of  $\mathbb{R}^2$  by the flows of  $X_{H^0}, X_{F^0}$ .



Each local involution  $\tau_{\pm}$  permutes these local orbits:

$$\tau_+(\mathcal{O}_1) = \mathcal{O}_3, \quad \tau_+(\mathcal{O}_2) = \mathcal{O}_4.$$

From topological point of view, an integrable system is a **singular Lagrangian foliation** of the phase space by integral manifolds  $\{H^\varepsilon = c_1, F^\varepsilon = c_2\}$ . General problem: describe topology of this foliation and its bifurcations along the above family of equilibria.

Questions about the singularity with hyperbolic 1:1 resonance:

- 1) To which **“standard form”** can be locally reduced the pair of functions  $H^\varepsilon, F^\varepsilon$  by some (non-linear, not necessarily symplectic) local coordinate changes?
  - 2) How do eigenvalues  $\lambda_1^\varepsilon$  and  $\lambda_2^\varepsilon$  bifurcate along the family and collide at  $\varepsilon = 0$ ?
  - 3) What are **local** topological invariants of (the bifurcation of) the foliation?
- What are **semi-global** topological invariants (in an invariant neighborhood of the equilibrium) of this foliation?
- 4) Is the given bifurcation (near the given equilibrium) **structurally stable**?

**Theorem 1.** Suppose that the integral manifold  $L^{0,0} = \{H^0 = F^0 = 0\}$  is compact, contains a 0-dimensional orbit (equilibrium) with non-semisimple hyperbolic resonance, and no other 0-dimensional orbits.

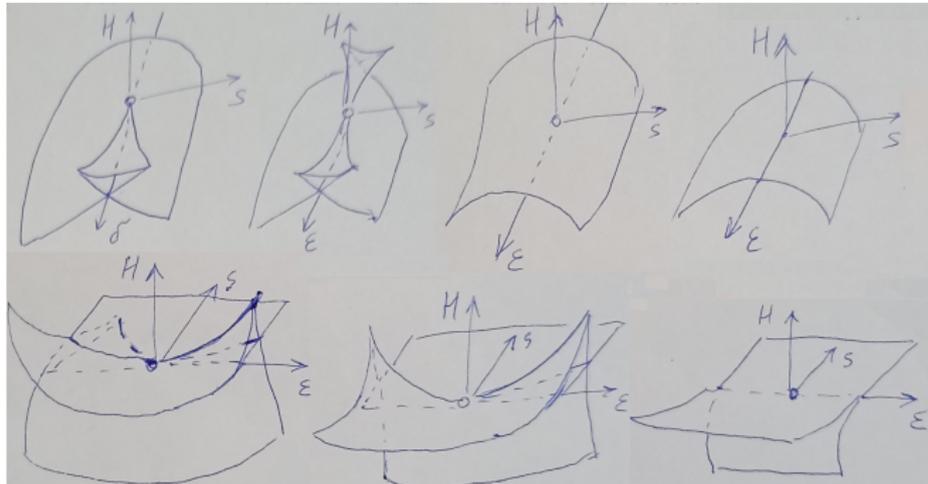
Then the pair of functions  $H^\varepsilon, F^\varepsilon$  can be **locally** reduced near this point, by some (non-linear, not necessarily symplectic) regular local coordinate changes, to one of the following standard forms:  $(\tilde{H}^\varepsilon, \tilde{F}^\varepsilon) = (G^{\delta(\varepsilon)}, S)$  where

$$G^\delta = \begin{cases} X + Y^2 + \delta Y, & \text{(local singularity } \mathcal{KL}_{1:1} \text{ [KL 2024])} \\ X + Y^3 + YS + \delta Y^2 + \frac{\delta^2}{4} Y, & \text{(new local singularity } \mathcal{KL}_{1:1}^\Delta) \end{cases}$$

$S = p_1 q_1 + p_2 q_2$ ,  $X = p_1 q_2$ ,  $Y = p_2 q_1$ ,  $Z = p_1 q_1 - p_2 q_2$ ,  $\delta$  is a detuning parameter,  $\delta(\varepsilon) = c\varepsilon^\ell$ .

The pair  $(c, \ell) \in \{(1, 2k-1), (1, 2k), (-1, 2k), (0, 0)\}$  is a smooth local invariant. It can be reduced to the case of  $k = 1$  via topological equivalences.

- ▶ Bifurcations  $\mathcal{KL}_{1:1}$  with  $\delta(\varepsilon) = \varepsilon$  are locally structurally stable. Bifurcation  $\mathcal{KL}_{1:1}^\Delta$  with  $\delta(\varepsilon) = \varepsilon$  is semi-globally structurally stable (but not locally).
- ▶ Local bifurcation diagram is a local topological invariant of the above singularities. They are shown in Figure.



Spectrum  $\text{spec } A^\varepsilon = \{\pm\lambda_1^\varepsilon, \pm\lambda_2^\varepsilon\}$  along the family of equilibria has the following behaviour near the 1:1 resonance:

► For  $\mathcal{KL}_{1:1}$ :

$(c, \ell) = (1, 1)$ ,  $\delta(\varepsilon) = \varepsilon$ ,  $\lambda_{1,2} = \lambda(\varepsilon) \pm \sqrt{\varepsilon}$  (splitting),

$(c, \ell) = (1, 2)$ ,  $\delta(\varepsilon) = \varepsilon^2$ ,  $\lambda_{1,2} = \lambda(\varepsilon) \pm \varepsilon$  (passing real),

$(c, \ell) = (-1, 2)$ ,  $\delta(\varepsilon) = -\varepsilon^2$ ,  $\lambda_{1,2} = \lambda(\varepsilon) \pm i\varepsilon$  (passing complex),

$(c, \ell) = (0, 0)$ ,  $\delta(\varepsilon) = 0$ ,  $\lambda_1 = \lambda_2 = \lambda(\varepsilon)$  (merging real).

► For  $\mathcal{KL}_{1:1}^\Delta$ :

$(c, \ell) = (1, 1)$ ,  $\delta(\varepsilon) = \varepsilon$ ,  $\lambda_{1,2} = \lambda(\varepsilon) \pm \varepsilon/2$  (passing real),

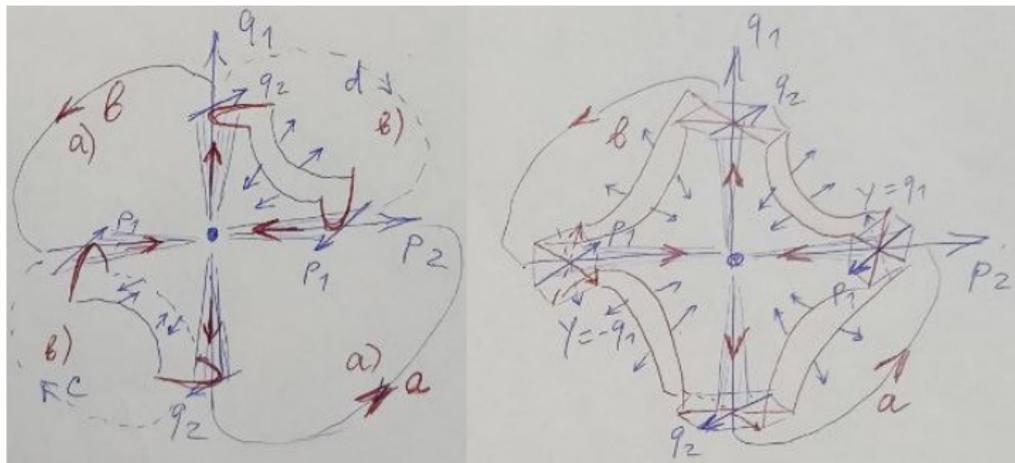
$(c, \ell) = (\pm 1, 2)$ ,  $\delta(\varepsilon) = \pm\varepsilon^2$ ,  $\lambda_{1,2} = \lambda(\varepsilon) \pm \varepsilon^2/2$  (passing real),

$(c, \ell) = (0, 0)$ ,  $\delta(\varepsilon) = 0$ ,  $\lambda_1 = \lambda_2 = \lambda(\varepsilon)$  (merging real).

**Theorem 2.** For each local singularity, there correspond **two semi-global** singularities (in a small invariant neighborhood of the equilibrium) determined by the following condition: a)  $\mathcal{O}_1 = \mathcal{O}_2$ , b)  $\mathcal{O}_1 = \mathcal{O}_3$ .

The integral manifold  $L^{0,0} = \{H^0 = F^0 = 0\}$  with its decomposition into orbits is a cell complex of the form  $L^{0,0} = ((a \vee b) \cup_W D^2) \cup_{W^{-1}} D^2$ , where the 2-cell is attached to the 1-skeleton in according to the word a)  $W = aba^{-1}b^{-1}$ , b)  $W = a^2b^{-2}$ , where  $a \approx b \approx S^1$ .

Topology of  $L^{0,0}$  is a semi-global topological invariant: a)  $\pi_1(L^{0,0}) = \pi_1(\mathbb{T}^2)$ , b)  $\pi_1(L^{0,0}) = \pi_1(KI^2)$ .



**Theorem 3.** Suppose that

$$H^\varepsilon = \lambda S + X + aY^2 + c\varepsilon^\ell Y + \dots,$$

where  $(c, \ell) \in \{(1, 2k-1), (1, 2k), (-1, 2k), (0, 0) \mid k \in \mathbb{N}\}$  as in Theorem 1.

If  $a \neq 0$  then the bifurcation has local type  $\mathcal{KL}_{1:1}$  with the corresponding pair  $(c, \ell)$ .

If  $a = 0$  then the bifurcation has local type  $\mathcal{KL}_{1:1}^\Delta$  with the corresponding pair  $(c, \ell)$ .

Thank you for your attention!